

## CORRIGENDUM

Application of the triple-deck theory of viscous–inviscid interaction  
 to bodies of revolution

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As indicated by Kluwick *et al.* in their paper (*J. Fluid Mech.* vol. 140, 1984, pp. 281–301), there is a mathematical error in (15) of our paper. This error, which originated from a typographical error in (8.3.8) of the book by Ward (*Linearized Theory of Steady High-Speed Flow*, C.U.P., 1955), affects the results presented in figures 3, 4 and 8. A few formulæ together with some typographical errors need to be corrected. In (14),  $\delta^{1/2}\epsilon^2$  and  $\delta^{1/2}/a$  should be replaced by  $\delta^{1/2}\epsilon^3$  and  $\delta^{1/2}/a$  respectively. A surplus left bracket in equation (19) should be deleted. In the denominator of (15),  $K_0(s)$  should be replaced by  $K_1(s)$ . The expression  $T'(z)$  in (19) should be replaced by  $T(z) = K_1(z)/K_0(z)$ . In the denominator of (24),  $2K_0(r_0\kappa_1) - K_1(r_0\kappa_1)^2$  should be replaced by  $K_0^2(r_0\kappa_1)$ . Equations (25)–(33) do not need any change. With these corrections, the function  $T'(z)$  has the same behaviour when  $z \rightarrow \infty$  as that previously given but the behaviour when  $z \rightarrow 0$  is different. As a consequence the integrands in (28) and (33) have the same approximate asymptotic expressions when  $\eta \rightarrow \infty$  as those given previously, but both become regular at the lower integration limit  $\eta = 0$ . Thus the numerical integration procedure described in the Appendix still stands except

$r_0$	0.5	1.0	5.0	10.0	$\infty$
$\kappa_1$	1.21439	1.06752	0.89259	0.86196	0.82716

TABLE 1. Zero of  $N(-i\kappa)$  as a function of  $r_0$

$r_0$	0.5	1.0	5.0	10.0	$\infty$
$x = 0.0$	0.35587	0.44673	0.63363	0.68316	0.75000
$x = 0.5$	0.36587	0.47621	0.70584	0.76725	0.85025
$x = 1.0$	0.32827	0.45134	0.71826	0.79151	0.89120
$x = 2.0$	0.24847	0.38014	0.70241	0.79795	0.93137
$x = 5.0$	0.11478	0.21815	0.59473	0.74086	0.97045
$x = 10.0$	0.05238	0.10898	0.43825	0.62389	0.98636
$x = 20.0$	0.02477	0.05108	0.26345	0.45123	0.99388

TABLE 2. Pressure distribution  $P(x)$

$r_0$	0.5	1.0	5.0	10.0	$\infty$
$x = 0.0$	-0.55564	-0.64007	-0.80575	-0.84875	-0.90654
$x = 0.5$	-0.14796	-0.22853	-0.40334	-0.45140	-0.51709
$x = 1.0$	-0.03897	-0.10739	-0.27948	-0.33066	-0.40259
$x = 2.0$	0.03581	-0.00501	-0.15586	-0.20969	-0.29064
$x = 5.0$	0.04225	0.04570	-0.02867	-0.07767	-0.17112
$x = 10.0$	0.01833	0.02893	0.01817	-0.01424	-0.10969
$x = 20.0$	0.00640	0.01127	0.02409	0.01406	-0.06936

TABLE 3. Shear-stress distribution  $(\tau - 1)/\alpha$

that the integration of (33) from 0 to  $A$ ,  $A$  to  $B$ , and from  $B$  to  $\infty$  should be replaced simply by that from 0 to  $B$  with the use of Simpson's rule and from  $B$  to  $\infty$  by analytical evaluation. Some typical values associated with the pressure and shear-stress distributions, calculated by the use of the corrected formulæ, are shown in the tables 1–3, which agree very well with those shown in figures 3 and 4 of the paper by Kluwick *et al.* (1984) found by a different approach.